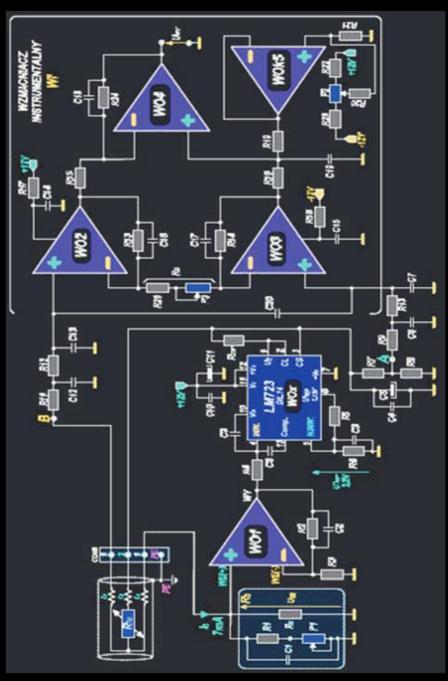


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PRZEGLĄD ELEKTROTECHNICZNY Vol 2022, No 9

Contents

01	Dirk Riedinger - A review and analysis of Quade's fundamental geometric time domain concept for the summation of non-	1
02	active powers of poly-phase systems Přemysl JANŮ, Šimon Brázda, Barbora ODVÁRKOVÁ - A module for analysis of power supplies of an electronic security systems	9
03	Inna HONCHARUK, Ihor KUPCHUK, Vitaly YAROPUD, Ruslan KRAVETS, Serhiy BURLAKA, Valerii HRANIAK, Julia POBEREZHETS, Volodymyr RUTKEVYCH - Mathematical modeling and creation of algorithms for analyzing the ranges of	14
04	the amplitude-frequency response of a vibrating rotary crusher in the software Mathcad Agnieszka CHOROSZUCHO - Analysis of the influence of the conductivity of bricks, type of bricks and size of the hollows on	21
05	the values of the electric field intensity Vadim MANUSOV, Dmitry ORLOV, Vitaly KARMANOV, Alexander KHUSNUTDINOV, Sergey KOKIN, Murodbek	26
06	SAFARALIEV - Analysis of electricity consumption forecasting methods for the coal industry Naima HASSANI, Hamid BOUZEBOUDJA, Mimoun YOUNES - Economic dispatch problem using artificial bee colony optimization based on predator and prey concept	32
07	Darmansyah, Imam ROBANDI - Intelligent Voltage Controller Based on Fuzzy Logic for DC-DC Boost Converter	36
80	M. I. Idris, H. Hussin, S. A. M. Junos, W. H. M. Saad, M. N. Shah Zainudin, N. M. Yatim, A. S. Jaafar, M. R. Kamarudin, K. M. Saipullah - VSLAM analysis using various ORBSLAM parameters setting	41
09	Benmessaoud Mourad, Abboun A.M, Benmessaoud N - Optimization of Direct Methanol Fuel Cell Power System	46
10	Arkadiusz WILCZYŃSKI, Olena HEBDA - Switch-mode power supply with a service life detection system based on monitoring of aluminium electrolytic capacitor parameters	51
11	Marian HYLA - Reactive power compensation in a 6 kV power grid supplying a 12-pulse thyristor hoisting machine	56
12	Adam GARCZAREK, Dorota STACHOWIAK – Static tests of current collectors in the production process and during	62
13	operation Bartłomiej STEFAŃCZAK, Dariusz ZIELIŃSKI - The problem of the influence of phase current asymmetry and the direction	67
	of energy flow on the idle current of a transformer in star-delta winding configuration	
14	Krzysztof WRÓBEL, Krzysztof TOMCZEWSKI, Andrzej TOMCZEWSKI - Analysis of the possibility of expanding a photovoltaic installation with a wind generator	73
15	Krzysztof SOŁTYS, Sebastian BARTEL, Krzysztof KLUSZCZYŃSKI - Installation for the production of ferromagnetic powder by electrolysis in laboratory conditions	77
16	Kalina DETKA, Sebastian LIGEZA - Investigation of the properties of an inductor with a hybrid core	82
17	Michał WYSOCKI, Robert NICPOŃ, Marta TRZASKA, Agnieszka CZAPIEWSKA - Research of Accuracy of RSSI	86
18	Fingerprint-Based Indoor Positioning BLE System Krzysztof GÓRECKI, Przemysław PTAK, Jakub HELENIAK - Investigations of the distortion of the supply current and optical	90
19	parameters of selected lamps Adrian PIETRUSZKA, Paweł GÓRECKI, Jacek TARASIUK, Agata SKWAREK - Influence of the location of solder the thermal parameters of the MOSFET transistors	94
20	Leszek PIECHOWSKI, Jan IWASZKIEWICZ, Adam MUC, Krystian KASPEREK - 4-channel temperature meter using Pt100 sensors	99
21	Damian BISEWSKI - Application of the genetic algorithm in the estimation process of models parameters of semiconductor devices	103
22 23	Krzysztof GÓRSKI, Michał ŁUKOMSKI - Renewable energy sources in powering miniature electronic devices Leszek PIECHOWSKI, Jan IWASZKIEWICZ, Adam MUC - The connection structure of the resistors is hidden in the value of a	107 111
24	natural binary code Jan IWASZKIEWICZ, Adam MUC, Agata BIELECKA - Six-phase Two-level VSI Control based on Polar Voltage Space Vectors	115
25	Magdalena BUDNAROWSKA, Jerzy MIZERACZYK - Simulation and experimental studies of the electromagnetic properties	119
26	of a planar metamaterial array structure in the microwave range Sergii BESPALKO, Jerzy MIZERACZYK - The plasma discharges in the anodic and cathodic regimes of plasma driven solution electrolysis for hydrogen production	122
27	Magdalena BUDNAROWSKA, Jerzy MIZERACZYK - Numerical simulation in the frequency domain of the shielding effectiveness of the interior of the enclosure with an opening against an ultrashort high power electromagnetic pulse	126
28	Krzysztof GÓRSKI, Sebastian SZYMAŃSKI, Igor MIELCZAREK, Jakub GRZESIAK - The literature review on detection and combating Unmanned Aerial Systems	130
29	Agnieszka CZAPIEWSKA, Ryszard STUDAŃSKI, Andrzej ŻAK - Data transmission in the hydroacoustic channel under	135
30	NLOS Kalina DETKA, Krzysztof GÓRECKI - Analysis the usefulness of the method of determining power losses in magnetic	139
31	materials based on the measurement of the hysteresis loop surface area Michał DOWNAR-ZAPOLSKI, Piotr MYSIAK - Analysis and synthesis, design and implementation of the single-phase voltage	143
32	inverter model controlled by PFM modulation Radoslaw GŁĄB, Dorota Rabczuk - Application to monitor the base station of a radio communication system using the SNMP	147
33	protocol Emilian ŚWITALSKI, Krzysztof GÓRECKI - Using Lambert-W function in modelling characteristics of an alkaline electrolyser Wiesław CITKO, Wiesław Sieńko - Image Recognition and Reconstruction as Inverse Problem, Using Machine Learning	150
34 35	System Krzysztof GÓRECKI, Paweł GÓRECKI - Influence of the form of the thermal model on accuracy of computing DC	154 158
	characteristics of IGBT module	
36 37	Dominika DĄBRÓWKA, Robert P. SARZAŁA - Suppression of the thermal crosstalk effect in laser arrays Robert P. SARZAŁA, Julita POBORSKA - Thermal Investigation of GaAs-Based 2D VCSEL Diode Arrays	162 166
38	Witalis PELLOWSKI Agnieszka IWAN , Krzysztof A. BOGDANOWICZ2 - Conversion of photons of light generated by radiation excited photoluminescence into electricity in izo-photovoltaic cells a new challenge for energy security	170
39	Jacek KONOPACKI, Jan MACHNIEWSKI - A simple, digital method for background estimation of timing mismatches in time-interleaved ADCs -	174

PRZEGLĄD ELEKTROTECHNICZNY Vol 2022, No 9

Contents

40	Agnieszka IWAN, Krzysztof A. BOGDANOWICZ, Wojciech PRZYBYŁ - Organic hole transporting layer in polymer and	178
41	perovskite solar cells – selected materials and technical aspects Wojciech PRZYBYŁ, Ireneusz PLEBANKIEWICZ, Krzysztof A. BOGDANOWICZ, Agnieszka IWAN - Detections methods	182
42	of structural defects in layers of organic solar cells Karolina KOWALSKA, Marta KUWIK, Joanna PISARSKA, Dominik DOROSZ, Wojciech A. PISARSKI - Multicomponent	186
-	titanate-germanate glasses for infrared photonics	
43 44	Jacek CHĘCIŃSKI, Zdzisław FILUS - Virus- and bactericidal LED lamp with increased efficiency Piotr MAĆKÓW, Piotr GUZDEK, Jacek BISKUPSKI, Wojciech GRZESIAK - Innovative energy storage with monitoring and	190 194
45	supervision functions. Selected Issues Jakub KISAŁA, Andrzej KOCIUBIŃSKI, Karolina CZARNACKA, Mateusz GĘCA, Jakub DUK - Investigations of the	199
46	influence of copper layer thickness on the giant magnetoresistance effect in NiFe/Cu/NiFe thin films Ewa SCHAB-BALCERZAK, Jolanta KONIECZKOWSKA, Karolina BUJAK - Controlling the photoinduced properties of	202
47	azopolymers Julian BALCEREK, Adam KONIECZKA, Paweł PAWŁOWSKI, Cyprian DANKOWSKI, Mateusz FIRLEJ, Piotr FULARA -	205
48	Automatic recognition of vehicle wheel parameters Adam KONIECZKA, Szymon BALAWAJDER - Laboratory stand for the automatic testing of digital cameras parameters	209
49	Małgorzata RUDNICKA, Ewa KLUGMANN-RADZIEMSKA - Consequences of suboptimal design of Building Integrated Photovoltaic Installation – a case study for Gdansk, Poland	213
50	Jakub SUDER - Parameters evaluation of cameras in embedded systems	216
51	Monika Pokladko-Kowar, Katarzyna Wojtasik - Organic light-emitting diodes based on pyrazoloquinoxaline derivatives	220
52	Paweł KIELAN, Szymon OLEŚ – An augmented reality application design dedicated to multi-module valve terminals in the Festo CPX series	224
53	Jarosław WRÓBEL, Sebastian ZŁOTNIK, Jacek BOGUSKI, Marek KOJDECKI, Jerzy WRÓBEL - Characterization of multi- channel charge carrier transport for epitaxial semiconductor structures	228
54	Katarzyna DYNDAŁ, Zbigniew SOBKOW, Jerzy SANETRA, Konstanty W. MARSZALEK - Hall effect test bench for	231
55	temperature dependence of carrier concentration Jan GIELZECKI, Ryszard MANIA, Konstanty MARSZALEK, Robert WOLANSKI - Deposition of Thin (Ti,Si)N Reflecive	235
56	Layers on Textiles Substrates S. IWANEK, Jerzy SANETRA, Konststny W. MARSZALEK - Optical method for homogeneity testing of thin films electrodes	239
	for photovoltaic cells	
57	Katarzyna DYNDAŁ, Konstanty W. MARSZALEK, Zbigniew KĄKOL - Design of complex energy systems (thin film photovoltaics, collectors, heat pumps and energy storage	243
58	Szymon KIEŁCZAWA, Artur WIATROWSKI - Specific method of deposition of aluminium-doped zinc oxide thin films on flexible glass substrates	247
59	Janusz RYBAK, Konstanty MARSZALEK - Deposition and optical properties investigations of WO₃ thin films for	251
60	electrochromic device applications Katarzyna UNGEHEUER, Konstanty W. MARSZALEK, Marzena MITURA-NOWAK, Zbigniew. KAKOL - Modification of	255
61	semiconducting copper oxide thin films using ion implantation Krzysztof Kluszczynski ,Zbigniew Pilch - Comparison of the magnetorheological 1- and 2-disc clutch in terms of overall	259
	dimensions, coefficients of the use of active materials, electric power consumption and spatial temperature distribution	
62 63	Maksymilian CIERNIEWSKI, Wojciech JAKUSZKO, Paweł STOBNICKI - Control of the main circuit of a hybrid rail vehicle Tomasz MAKOWSKI, Maciej GIBAS, Kacper CHOJKOWSKI - Design and construction of a wheelchair with a caterpillar	265 269
64	drive with a seat auto-leveling system Ewa ZAWADZKA, Henryk BRZEZIŃSKI, Urszula ŚWITAŁA - Secondary raw materials in composite products in the	273
65	automotive industry Milena KURZAWA, Rafał M. WOJCIECHOWSKI - Determination of the parameters of equivalent schema of the pulse	277
	transformer using the Cauer circuits	
66 67	Jakub BERNATT, Stanisław GAWRON, ^T adeusz GLINKA, Artur POLAK - Block Transformer Fires Marek HRECZKA, Wojciech BURLIKOWSKI, Janusz HETMAŃCZYK, Aleksandra Kolano-BURIAN, Roman KOLANO -	282 126
68	Core of brushless DC motors made of amorphous soft magnetic material Maciej JAROSZEWSKI, Paweł RÓZGA, Abderrahmane BEROUAL, Karol SKOWROŃSKI, Mariusz WÓJCIAK - Some	293
	factors that influence the breakdown voltage of biodegradable transformer oil	
69	Bogusław GRUSZCZYŃSKI, Rafał M. WOJCIECHOWSKI - Design and synthesis of the magnetic circuit of a differential eddy current transducer for non-destructive testing	296
70	Jakub BERNATT, Stanisław GAWRON, Tadeusz GLINKA, Tomasz WOLNIK - Analysis of the benefits of using 6-phase windings in motors excited with permanent magnets	302
71	Krzysztof Mars, Mateusz Sałęga – Starzecki, Elżbieta Godlewska - Influence of the magnetron power mode on the properties of SnSe layers	307
72	Sławomir Krawczyk, Elżbieta Czerwosz, Halina Wronka, Piotr Firek, Mariusz Sochacki, Jan Szmidt - Comparison of Field Effect Transistor with C-nPd gate and resistive C-nPd film sensing properties toward hydrogen	311

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Mathematical modeling and creation of algorithms for analyzing the ranges of the amplitude-frequency response of a vibrating rotary crusher in the software Mathcad

Abstract. The article is devoted to the study of motion laws for rotary vibration crusher. Kinematic and dynamic analysis was performed. Differential equations of rotor motion are solved and analyzed, frequency response and energy consumption graphs in MathCad 15.0 software environment are presented. Verification of the mathematical model was carried out by comparing the results of experimental research with theoretical research. It was proved that the proposed mathematical models are adequate (discrepancy are 7.2 to 12.1%).

Streszczenie. Artykuł poświęcony jest badaniu praw ruchu obrotowego kruszarki wibracyjnej. Przeprowadzono analizę kinematyczną i dynamiczną. Równania różniczkowe ruchu wirnika są rozwiązywane i analizowane, prezentowane są wykresy odpowiedzi częstotliwościowej i zużycia energii w środowisku oprogramowania MathCad 15.0. Weryfikację modelu matematycznego przeprowadzono poprzez porównanie wyników badań eksperymentalnych z badaniami teoretycznymi. Wykazano, że zaproponowane modele matematyczne są adekwatne (rozbieżność wynosi od 7,2 do 12,1%). (Modelowanie matematyczne i tworzenie algorytmów analizy zakresów odpowiedzi amplitudowo-częstotliwościowej wibracyjnej kruszarki obrotowej w programie Mathcad)

Keywords: Lagrange's equations II kind's, differential equations, parameters of oscillations, analytical mechanics, vibrating machine **Słowa kluczowe:** równanie różnicowe, odpowiedź częstotliwościowa, kruszarka wibracyjna

Introduction

For feeding livestock very often use cheap fodder grain, which is not suitable for food purposes. Usually such grain has high humidity [1, 2]. Dry food grain is rarely used. But the specific energy consumption, which is specified in the technical characteristics for hammer crushers are reliable only for quality dry grain (humidity about 14-16%) [2, 3]. During the grinding of wet (fodder) grain, productivity is significantly reduced and specific energy consumption increases [4]. Therefore, there is a need to develop an equally energy-efficient crusher for both dry and wet grain.

Analysis of literary sources and problem statement

During the destruction process of materials with plasticity properties, energy is used to overcome molecular bonds and irreversible plastic deformations [5, 6]. The energy spent on deformation is converted into heat [7, 8, 9].

With increasing moisture content, fragility and ultimate strength decrease, ductility and ultimate deformation before fracture increase [5, 10].

In addition, during the grinding process of forage grains can often be adhesion to the sieve of the wet shredded product [10, 11, 12].

The development of a vibratory crusher consists of many complex stages [13, 14]. The first steps in the development process of a machine are theoretical studies [15, 16]:

- 1) development of kinematic scheme (the new design should solve the problem of effective grinding (dry and wet grain) and prevent blockage of the sieve with crushed material);
- 2) kinematics research (these results are the basis for dynamic analysis) [17];
- 3) dynamics research (these results are the basis for calculating the crusher drive and creating a design drawing) [18].

Using the infrastructure of the laboratory of the process and processing of equipment and food industry Vinnytsia National Agrarian University, a design of vibratory crusher was developed [6, 19, 20]. Sharp disks will be installed in

this crusher instead of rectangular plates (hammers). That is the methods of impact and cutting to destroy the material will be combined [19]. Thus, a local overvoltage of surface micro volumes at the places of application of loads will be created [13, 21]. When the disc hits the grain quickly, it causes the sharp blade to sink into the body and create the pressure needed to break the material [8, 22]. In addition, for more intensive sieving of the product that already crushed oscillations of the sieve will be provided [19, 23].

Purpose and tasks of research

The purpose of the research is to development of a mathematical model for the crusher being designed and determination of ranges of amplitude-frequency characteristics in which energy consumption will be rational.

To achieve this goal, it is necessary: a kinematics study performed and to obtain the laws of motion; develop differential equations describing the motion of the disk rotor and energy consumption by the drive; find the solutions of the differential equations and determine the most optimal ranges of values for amplitude-frequency characteristics; check whether mathematical models are adequate by comparing the results of theoretical and experimental studies using a laboratory model of a vibrating crusher.

Materials and methods

Scientific articles position based on the classical theory of mechanical oscillations of the laws of theoretical mechanics and physics, kinematics analysis was done analytical method, the principle of superposition [12, 15].

The machine is represented mathematical model with 6 degrees of freedom, namely the shifting of the centre of mass of the container along with the axis OX (Fig. 1), shifting the centre of mass of the rotor along with the axis OX, shifting the centre of mass of the container along the axis OZ, shifting the centre of mass of the rotor along the axis OZ, angular shifting of the rotor relative to the axis O1Y1, angular shifting of the disc relative to axis O2Y2 [24].

To determine the crusher's laws of motion along each of the independent coordinates $(x, y, z, \varphi_1, \varphi_2, \varphi_3)$, Lagrange's

equations of the second kind will be used [17, 24] (1). In this vibrational system (Fig. 1), four characteristic moving masses can be distinguished in the total mass m (2).

To determine the equations of the generalized linear velocity of the centres of mass of the structural elements of the vibratory disk crusher, the mechanism will be divided into elementary components (Fig. 2), which will be studied separately [25]. To solve and analyze the obtained equations of motion of the executive body of the vibratory disk crusher, the mathematical environment MathCad 15.0 was used [6]. The use of this program allowed us to determine the patterns of change of oscillation parameters depending on the angular velocity of the drive shaft.

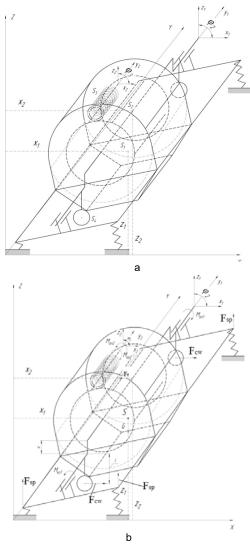


Fig. 1. Calculation scheme for studying: a) kinematics; b) forces and moments

(1)
$$\begin{cases} \frac{d}{dt} \left(\frac{\partial I}{\partial \dot{x}} \right) - \frac{\partial I}{\partial x} = Qx \\ \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{z}} \right) - \frac{\partial T}{\partial z} = Qz \\ \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{y}} \right) - \frac{\partial T}{\partial y} = Qy \\ \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{\phi}_{1}} \right) - \frac{\partial T}{\partial \phi_{1}} = Q\phi_{1} \\ \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{\phi}_{2}} \right) - \frac{\partial T}{\partial \phi_{2}} = Q\phi_{2} \\ \frac{d}{dt} \left(\frac{\partial T}{\partial \phi_{2}} \right) - \frac{\partial T}{\partial \phi_{2}} = Q\phi_{3} \end{cases}$$

where T – kinetic energy of the system; Qx, Qz, Qy, $Q\phi_1$, $Q\phi_2$, $Q\phi_3$ – generalized resistance forces.

$$\begin{cases} m = m_1 + m_2 + m_3 + m_4; \\ m_1 = m_{\kappa} + m_m + m_{sf} + m_b; \\ m_2 = m_r + m_c; \\ m_3 = m_d; \\ m_4 = m_{cw}; \\ m_r = m_{esh} + m_{cd} + m_{var} + m_{sup} + m_{axles}; \end{cases}$$

where m_{κ^-} mass of frame, kg; m_m^- mass of material, kg; m_{sf}^- mass of support frame, kg; m_b^- mass of bearing units, kg; m_r^- rotor weight, kg; m_c^- mass of couplings, kg; m_d^- mass of impact discs, kg; m_{cw}^- weight of counterweight, kg; m_{esh}^- mass of eccentric shaft, kg; m_{cd}^- mass of intermediate discs, kg; m_{var}^- a mass of eccentric variation mechanisms, kg; m_{sup}^- mass of support discs, kg; m_{axles}^- mass of disc axles, kg.

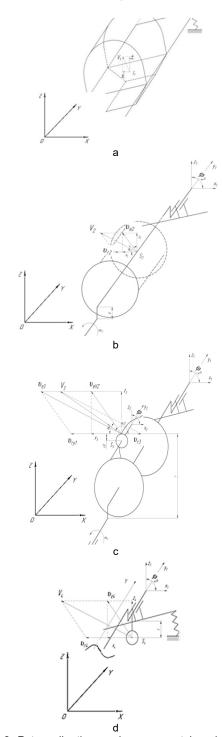


Fig. 2. Rotary vibration crusher: a - container; b - rotor; c - disc; d - counterweight.

As a result of research studies performed earlier, the approximate ranges of oscillation amplitude, vibration velocity, vibration acceleration and oscillation intensity were determined. However, to assess energy efficiency, it is necessary to experimentally investigate the effect of amplitude-frequency characteristics on energy consumption for the crusher drive.

Alsow, based on previous study [8, 20], a database was adopted, which included the values: the range of angular velocity of the drive shaft $\omega_2 = 0...150$ rad s⁻¹, and the time factor interval t = 0...60 s, as well as the values of the accepted constants [19] of the studied system (Table 1).

Table 1. Numerical values of the main constants adopted for the studied system

ioi ilie studied system	
Constants	Value
Total moving weight, kg	83.2
- m ₁ , kg	32.7
- m ₂ , kg	31.9
- m ₃ , kg	13.1
- m_4 , kg	5.5
The distance from the axis of rotation to the center of mass of the rotor e, m	0.005
Working disk radius r_d , m	0.045
The radius of the support disk r_{sup} , m	0.14
The distance from the top of the working disk to the axis of rotation r , m	0.19
The distance from the axis of rotation to the center of mass of the counterweight <i>I</i> , m	0.044
Stiffness of elastic elements C, N/m	
- along the axis <i>OX</i> : <i>C</i> _x	3900
- along the axis OZ : C _z	3900

The experimental study of ranges of amplitudefrequency characteristics at which energy consumption will be rational was carried out at Vinnytsia National Agrarian University using the laboratory model of a rotary vibration crusher [6].

To record the angular velocity values of the drive shaft, the UNI-T UT372 wireless tachometer (Uni-Trend Technology Limited, Dongguan, China) was used. To manage and change rotation frequencies of the motor shaft (up 0 to 150 s⁻¹ with step 5 s⁻¹), the AOSN-20-220-75 (PJSC «Megommeter», Uman, Ukraine) autotransformer was used [26]. To determine the energy consumption to drive the crusher, the EMF-1 (ELTIS Electric, Lviv, Ukraine) electronic wattmeter was used [20]. The operating principle and operating rules for thos devices are described in the technical documentation. To record the amplitudefrequency characteristics of the vibratory disk crusher, a sensor based on the ST Microelectronics accelerometer was developed [8, 27].

Taking into account the permissible errors of the measuring equipment [28], a critical value of the discrepancy between the experimental and theoretical research was taken by 15% [20]. Exceeding this boundary indicates the unreliability of the mathematical model and the inability to use it when designing a crusher of this type. Processing and analysis of the research results were carried out in the Microsoft Excel 2019 software environment.

Research results

To determine the linear velocity of the centres of mass of structural elements of oscillatory system we will divide the mechanism into elementary components are links and we will study them separately [15, 29]. Velocity (V₁, V₂, V₃, $V_4)$ for the generalized mass S_1 (container),), S_2 (rotor), S_3 (disc), S₄ (counterweight) (Fig. 2), m·s⁻¹:

(3)
$$V_1 = \sqrt{\dot{x}_1^2 + \dot{z}_1^2};$$

(4)
$$V_2 = \sqrt{v_{r2}^2 + v_{e2}^2 + 2v_{r2}v_{e2}\cos\psi_1};$$

(5)
$$V_3 = \sqrt{v_{r3}^2 + v_{e3}^2 + 2v_{r3}v_{e3}\cos\psi_3} ;$$

(6)
$$V_4 = \sqrt{v_{r4}^2 + v_{e4}^2 + 2v_{r4}v_{e4}\cos\psi_4}$$

where \dot{x}_1, \dot{z}_1 – the velocity S₁ along with the axis *OX* and *OZ*, m·s⁻¹; ν_{r2} , ν_{r3} , ν_{r4} – respectively the relative velocity of S_2 (the moving coordinate system – $x_1y_1z_1$), S_3 (the moving coordinate system – $x_2y_2z_2$) and S₄ (the moving coordinate system – $x_1y_1z_1$), m·s⁻¹; ν_{e2} , ν_{e3} , ν_{e4} , respectively the frame velocity of S₂ and coordinate system – $x_1y_1z_1$ (the fixed coordinate system - XYZ), S₃ and coordinate system $x_2y_2z_2$ (the fixed coordinate system - $x_1y_1z_1$), S₄ and coordinate system - x₁y₁z₁ (the fixed coordinate system -XYZ), m·s⁻¹; ψ_1 , ψ_3 , ψ_4 – respectively the angle between the vectors \vec{v}_{r2} and \vec{v}_{e2} , \vec{v}_{r3} and \vec{v}_{e3} , \vec{v}_{r4} and \vec{v}_{e4} , rad.

(7)
$$v_{r2} = e \cdot \dot{\varphi}_1, \ v_{e2} = \sqrt{x_2^2 + z_2^2} \cdot \dot{\varphi}_2 \rightarrow \sqrt{x_1^2 + z_1^2} \cdot \dot{\varphi}_2;$$

(8)
$$v_{r3} = r_o \cdot \dot{\varphi}_3 \cdot ku , v_{e3} = \sqrt{v_{ry1}^2 + v_{eXZ}^2 + 2v_{ry1}v_{eXZ}\cos\psi_2}$$

(9)
$$v_{r_4} = l \cdot \dot{\varphi}_2, v_{e_4} = \sqrt{x_4^2 + z_4^2} \cdot \dot{\varphi}_2 \rightarrow \sqrt{x_1^2 + z_1^2} \cdot \dot{\varphi}_2,$$

where: e - distance from S₂ to the axis 0_1y_1 , m; $\dot{\varphi}_2$ - the angular velocity of the rotor, rad $\cdot s^{-1}$; x_2 , z_2 – displacement of the S_2 relative to fixed axes OX and OZ, m; r_d – the radius crusher disc, m; $\dot{\phi}_3$ – the angular velocity crusher disc, rad·s⁻¹; v_{rv1} - the relative velocity of cutting edge (the moving coordinate system – $x_1y_1z_1$), m·s⁻¹; v_{exz} – the frame velocity of the cutting edge and coordinate system $-x_2y_2z_2$ (the fixed coordinate system – $x_1y_1z_1$), $m \cdot s^{-1}$; ψ_2 – the angle between the vectors \vec{v}_{eXZ} and \vec{v}_{ry1} , rad; I – distance from S₄ to the axis 0_1y_1 , m; x_4, z_4 – displacement of the S_4 relative to fixed axes OX and OZ, m; ku - the transfer coefficient of torque ($ku = 0 \dots 1$).

When rotating the working rotor equipment between the edge crusher disc and material, the friction forces F_{fm} , the rotation is working disk is only possible where $F_{fm} > F_{fa}$. F_{fa} is the friction forces in conjunction with friction «crusher disk – disk axis» [12]. If there is $ku \to 1$ an increase F_{fm} , whereas at ku=0, $F_{fm}=0$, and as a result $\dot{\varphi}_3$ =0 [12, 24]. If ku > 0, when

$$\dot{\varphi}_3 = \frac{\dot{\varphi}_2 \cdot r_{sup}}{r_s \cdot ku}$$

$$v_{\rm rel} = r \cdot \dot{\varphi}_2;$$

(10)
$$\phi_{3} = \frac{\dot{\varphi}_{2} \cdot r_{sup}}{r_{d} \cdot ku};$$
(11)
$$v_{ry1} = r \cdot \dot{\varphi}_{2};$$
(12)
$$v_{exz} = \sqrt{x_{3}^{2} + x_{3}^{2}} \cdot \dot{\varphi}_{2} \rightarrow \sqrt{x_{1}^{2} + z_{1}^{2}} \cdot \dot{\varphi}_{2}$$

where: r_{sup} - the radius of the support disk, m; r - the distance from the edge of the crusher disc to the axis 0_1y_1 , m; x_3 , z_3 – displacement of the S₃ relative to fixed axes OX and OZ, m.

Equation (3-6) for velocites V₁₋₄, takes the form:

$$(13) V_1 = \sqrt{\dot{x}_1^2 + \dot{z}_1^2};$$

(14)
$$V_2 = \sqrt{(e \cdot \dot{\varphi}_2)^2 + (x_1^2 + z_1^2)\dot{\varphi}_2 + 2 \cdot e \cdot \dot{\varphi}_2 \cdot x_1};$$

(13)
$$V_{1} = \sqrt{\dot{x}_{1}^{2} + \dot{z}_{1}^{2}};$$
(14)
$$V_{2} = \sqrt{(e \cdot \dot{\varphi}_{2})^{2} + (x_{1}^{2} + z_{1}^{2})\dot{\varphi}_{2} + 2 \cdot e \cdot \dot{\varphi}_{2} \cdot x_{1}};$$
(15)
$$V_{3} = \sqrt{(r_{d} \cdot \dot{\varphi}_{3} \cdot ku)^{2} + (r \cdot \dot{\varphi}_{2})^{2} + (x_{1}^{2} + z_{1}^{2}) \cdot \dot{\varphi}_{2} + 2(r \cdot \dot{\varphi}_{2})x_{1} - 2(r_{d} \cdot \dot{\varphi}_{3} \cdot ku)(2x_{1} + r \cdot \dot{\varphi}_{2})}$$

(16)
$$V_4 = \sqrt{(l \cdot \dot{\varphi}_2)^2 + (x_1^2 + z_1^2) \cdot \dot{\varphi}_2 + 2 \cdot l \cdot \dot{\varphi}_2 \cdot x_1}$$

Total kinetic energy of the system:

(17)
$$T = \frac{1}{2}m_{1}[\dot{x}_{1}^{2} + \dot{z}_{1}^{2}] + \frac{1}{2}m_{2}[(e \cdot \dot{\varphi}_{2})^{2} + (x_{1}^{2} + z_{1}^{2}) \cdot \dot{\varphi}_{2} + 2 \cdot e \cdot \dot{\varphi}_{2} \cdot x_{1}] + \frac{1}{2}m_{3}[(r_{d} \cdot \dot{\varphi}_{3} \cdot ku)^{2} + (r \cdot \dot{\varphi}_{2})^{2} + (x_{1}^{2} + z_{1}^{2}) \cdot \dot{\varphi}_{2} + 2(r \cdot \dot{\varphi}_{2})x_{1} - 2(r_{d} \cdot \dot{\varphi}_{3} \cdot ku) \times$$

$$\begin{split} &\times (2x_1 + r \cdot \dot{\varphi}_2)] + \frac{1}{2} m_4 [(l \cdot \dot{\varphi}_2)^2 + \\ &(x_1^2 + z_1^2) \cdot \dot{\varphi}_2 + 2 \cdot l \cdot \dot{\varphi}_2 \cdot x_1] + \frac{1}{2} [(l_2 \dot{\varphi}_2^2) \\ &+ (l_3 \dot{\varphi}_3^2 k u) + (l_3 \dot{\varphi}_2^2) + (l_4 \dot{\varphi}_2^2)]. \end{split}$$

The generalized force can be interpreted as a coefficient before the variation of the generalized coordinate in the expression for the sum of the elementary works of all active forces.

Using the calculation scheme in Fig. 1b generalized forces can be identified:

(18)
$$\begin{cases} Qx = (m_2 + m_3)\omega_2^2 e \cos(\omega_2 \cdot t) - m_4\omega_2^2 l \times \\ \times \cos(\omega_2 \cdot t) - c_x x \end{cases} \\ Qz = \begin{bmatrix} (m_2 + m_3)\omega_2^2 e \cdot \sin(\omega_2 \cdot t) - m_4\omega_2^2 l \times \\ \times \sin(\omega_2 \cdot t) - (m_1 + m_2 + m_3 + m_4)g - c_z z \end{bmatrix} \\ Q\dot{\varphi}_2 = \begin{bmatrix} M_{\text{KP1}} + m_2\omega_2^2 e^2 \cdot \sin(\omega_2 \cdot t) - m_4\omega_2^2 l^2 \times \\ \times \sin(\omega_2 \cdot t) - M_{\text{On1}} \end{bmatrix} \\ Q\dot{\varphi}_3 = (M_{\text{KP2}} - M_{\text{On2}}) \cdot ku \end{cases}$$
where $c_z = c_z$ the stiffness of the elastic elements along the

where c_x, c_z – the stiffness of the elastic elements along the respective axes.

Using the MathCad 15.0 software environment, first and second parts for expressions of equation (1) were solved analytically. Therefore, expressions of equation (1) will take the following form:

$$\ddot{x} + \alpha_{x} \cdot \dot{x} - x \left[\frac{\dot{\varphi}_{2}(m_{2} + m_{3} + m_{4}) + c_{x}}{m_{1}} \right] = cos(\omega_{2} \cdot t) \times$$

$$\times \left(\frac{(m_{2} + m_{3})\omega_{2}^{2}e - m_{4}\omega_{2}^{2}l}{m_{1}} + \frac{m_{2} \cdot \dot{\varphi}_{2}^{2} \cdot e}{m_{1}} + \frac{m_{3} \cdot r \cdot \dot{\varphi}_{2}^{2} - m_{3} \cdot 2ku \cdot r_{d} \cdot \dot{\varphi}_{3}^{2} + m_{4} \cdot \dot{\varphi}_{2}^{2} \cdot l}{m_{1}};$$

$$\ddot{z} + \alpha_{z} \cdot \dot{z} - z \cdot \left[\frac{\dot{\varphi}_{2}(m_{2} + m_{3} + m_{3}) + c_{z}}{m_{1}} \right] =$$

$$= \frac{(m_{2} + m_{3})\omega_{2}^{2}e - m_{4}\omega_{2}^{2}l}{m_{1}} \times sin(\omega_{2} \cdot t) - \frac{(m_{1} + m_{2} + m_{3} + m_{4})g}{m_{1}};$$

$$\frac{m_{3}(2rx - 2\ddot{\varphi}_{2}r^{2} - 2ku \cdot r_{d}\dot{\varphi}_{3} \cdot r + x^{2} + z^{2})}{m_{1}} + \frac{m_{2}(x^{2} + 2ex + z^{2} + 2\ddot{\varphi}_{2}e^{2})}{2} + \frac{m_{4}(2\ddot{\varphi}_{2}l^{2} + 2lx + x^{2} + z^{2})}{2} + \frac{+ \dot{\varphi}_{2}(l_{2} + l_{3} + l_{4}) = M_{\text{kp1}} + m_{2}\omega_{2}^{2}e^{2} \cdot sin(\omega_{2} \cdot t) - m_{4}\omega_{2}^{2}l^{2} \cdot sin(\omega_{2} \cdot t) - m_{0n1};$$

$$(22) \qquad ku \cdot r_{d} \cdot m_{3}(ku \cdot r_{d} \cdot \ddot{\varphi}_{3} - (2x + r \cdot \dot{\varphi}_{2})) + l_{3} \cdot ku \cdot \ddot{\varphi}_{3} = \\ = (M_{\text{KD2}} - M_{\text{On2}}) \cdot ku.$$

To establish the laws of motion for the axes OX and OZ, it is sufficient to use only the first two equations (19, 20), which after mathematical transformations will take the form:

$$(23) \quad \alpha_{x} \cdot \dot{x} - x \left[\frac{\dot{\varphi}_{2}(m_{2} + m_{3} + m_{4}) + c_{x}}{m_{1}} \right] = cos(\omega_{2} \cdot t) \times \\ \times \left[\left(\frac{(m_{2} + m_{3})\dot{\varphi}_{2}^{2}e - m_{4}\dot{\varphi}_{2}^{2}l}{m_{1}} \right) + \frac{1}{m_{2} \cdot \dot{\varphi}_{2}^{2} \cdot e + m_{3} \cdot r \cdot \dot{\varphi}_{2}^{2} - m_{3} \cdot 2ku \cdot r_{d} \cdot \dot{\varphi}_{3}^{2} + m_{4} \cdot \dot{\varphi}_{2}^{2} \cdot l}{cos(\omega_{2} \cdot t)m_{1}} \right]; \\ \ddot{z} + \alpha_{z} \cdot \dot{z} - z \cdot \left[\frac{\dot{\varphi}_{2}(m_{2} + m_{3} + m_{4}) + c_{z}}{m_{1}} \right] = sin(\omega_{2} \cdot t) \times \\ \times \left[\frac{(m_{2} + m_{3})\omega_{2}^{2}e - m_{4}\omega_{2}^{2}l}{m_{1}} - \frac{(m_{1} + m_{2} + m_{3} + m_{4})g}{sin(\omega_{2} \cdot t) \cdot m_{1}} \right].$$

The solutions of equations (23) and (24) will be found for both second-order linear differential equations with constant coefficients, assuming that $\dot{\varphi}_2 = \omega_2$. The dissipation coefficients of this system can be represented as [17]:

 $\alpha_x = 2\sqrt{k_x^2 - \omega_2^2}; \ \alpha_z = 2\sqrt{k_z^2 - \omega_2^2}.$ (25)Specific modulus of forcing force (Lanets et al., 2019):

(26)
$$F_{mx} = \left(\frac{(m_2 + m_3)\omega_2^2 e - m_4\omega_2^2 l}{m_1}\right) + \frac{m_2 \cdot \omega_2^2 \cdot e + m_3 \cdot r \cdot \omega_2^2 - m_3 \cdot 2ku \cdot r_d \cdot \omega_3^2 + m_4 \cdot \omega_2^2 \cdot l}{cos(\omega_2 \cdot t)m_1} + \frac{(m_2 + m_3)\omega_2^2 e - m_4\omega_2^2 l}{m_1} - \frac{(m_1 + m_2 + m_3 + m_4)g}{sin(\omega_2 \cdot t) \cdot m_1}.$$

The natural frequency of oscillations of the system relative to the axis OX and OZ [24]:

(28)
$$k_x^2 = \frac{\omega_2(m_2 + m_3 + m_4) + c_x}{m_1};$$
(29)
$$k_z^2 = \frac{\omega_2(m_2 + m_3 + m_4) + c_z}{m_1}.$$

By solving the obtained equations as linear differential equations of the second order with constant coefficients, the dependences of the motion of the executive body of the studied machine are obtained. Due to the scattering of energy in the system under study, the free oscillations are damped, as a result of which the obtained equations for the steady-state will take the form [17, 23]:

(30)
$$x = \frac{F_{mx}\alpha_{x}\omega_{2}}{(k_{x}^{2} - \omega_{2}^{2})^{2} + \alpha_{x}^{2}\omega_{2}^{2}} \sin(\omega_{2}t) + \frac{F_{mx}(\omega_{2}^{2} - k_{x}^{2})}{(k_{x}^{2} - \omega_{2}^{2})^{2} + \alpha_{x}^{2}\omega_{2}^{2}} \cos(\omega_{2}t);$$

$$z = \frac{F_{mz}(k_{z}^{2} - \omega_{2}^{2}) \sin\omega_{2}t}{(k_{z}^{2} - \omega_{2}^{2})^{2} + \alpha_{z}^{2}\omega_{2}^{2}} + \frac{F_{mz}\alpha_{z}\omega_{z}\cos\omega_{2}t}{(k_{z}^{2} - \omega_{2}^{2})^{2} + \alpha_{z}^{2}\omega_{2}^{2}}.$$

The amplitude of oscillations about the axis OX and OZ has the form:

$$A_{\chi} = \frac{F_{mx}}{\sqrt{(k_x^2 - \omega_z^2)^2 + \alpha_x^2 \omega_z^2}};$$
(33)
$$A_{Z} = \frac{F_{mz}}{\sqrt{(k_z^2 - \omega_z^2)^2 + \alpha_z^2 \omega_z^2}}.$$
The absolute amplitude of oscillations:

(34)
$$A = \sqrt{A_x^2 + A_z^2}.$$

Taking into account the equation (32) and (33):

$$(35) \quad A = \begin{pmatrix} \frac{\left(\frac{(m_2 + m_3)\omega_2^2 e - m_4\omega_2^2 l}{m_1}\right) + \frac{m_2 \cdot \omega_2^2 \cdot e + m_3 \cdot r \cdot \omega_2^2 - m_3 \cdot 2ku \cdot r_d \cdot \omega_3^2 + m_4 \cdot \omega_2^2 \cdot l}{cos(\omega_2 \cdot t)m_1} \\ + \frac{\left(\frac{\omega_2(m_2 + m_3 + m_4) + c_x}{m_1} - \omega_2^2\right)^2 + \frac{m_2}{m_1}}{\left(\frac{\omega_2(m_2 + m_3 + m_4) + c_x}{m_1} - \omega_2^2\right)^2 \cdot \omega_2^2} \end{pmatrix} + \frac{\left(\frac{(m_2 + m_3)\omega_2^2 e - m_4\omega_2^2 l}{m_1} - \frac{m_1}{c(m_1 + m_2 + m_3 + m_4)g} - \frac{m_1}{sin(\omega_2 \cdot t) \cdot m_1}\right)^2}{\left(\frac{(\omega_2(m_2 + m_3 + m_4) + c_x}{m_1} - \omega_2^2\right)^2 + \frac{m_1}{m_1}\right)^2} + \frac{\left(\frac{(\omega_2(m_2 + m_3 + m_4) + c_x}{m_1} - \omega_2^2\right)^2 + \frac{m_1}{m_1}\right)^2}{\left(\frac{(\omega_2(m_2 + m_3 + m_4) + c_x}{m_1} - \omega_2^2\right)^2 + \frac{m_1}{m_1}\right)^2}$$

Also, using equation (35), the parameters of the vibration field for the steady-state, which are proportional to the amplitude and frequency of oscillations, namely: vibration velocity $u = A \cdot \omega$; vibration acceleration $a = A \cdot \omega^2$; oscillation intensity $I = a \cdot v = A^2 \cdot \omega^3$ [15].

The power of the drive of the studied machine [19]:

(36)
$$N_{pd} = (N_{Fmax} + N_{fr})/\gamma_{dr}$$
, where: N_{Fmax} – maximum power developed by the forcing force to generate the required oscillations; N_{fr} – power consumption for friction in the support nodes; γ_{dr} – efficiency of the drive.

Power developed by the forcing force:

 $(37) N_F = F_m \cdot v,$

Power consumption for friction [20]:

$$N_{fr} = 0.5 \cdot F \cdot \mu \cdot d_{sh} \cdot \omega_2^2,$$

where: $\mu = 0.05...0.08$ – coefficient of friction; d_{sh} – the diameter of the drive shaft ($d_{sh} = 0.04 \, m$).

For further analysis, we used the numerical values of constants and other experimental data obtained as a result of exploratory research on the basis in laboratory Vinnytsia National Agrarian University. The software algorithm that was created to analyze the amplitude-frequency ranges of the crusher and estimate the energy consumption of the drive is based on developed mathematical model, which is presented in the form of differential equations (17-38), which were entered into the working field of the software environment MathCad 15.0 [6]. Visualization of the obtained results is carried out in the form of an array of numerical values, three-dimensional graphs of amplitude-frequency and energy characteristics. In fig. 3 shows a fragment of the listing for automated determination of influence of design and technological parameters on the values amplitude of vibrations, vibration speed, vibration acceleration, vibration intensity and power consumption of the drive.

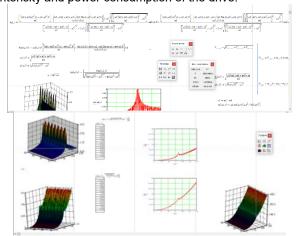


Fig. 3. Fragment of software algorithm for analyze the amplitude-frequency ranges of the vibratory crusher and visualization of the obtained results

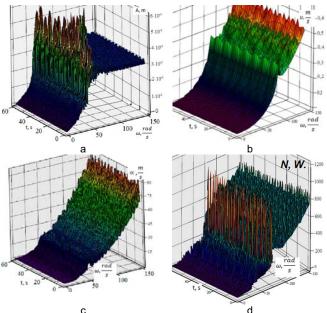


Fig. 4. Amplitude-frequency and energy characteristics of vibrating crusher: a) amplitude of oscillations; b) vibration speed; c) vibration acceleration; d) power consumption

As a result of mathematical analysis of compound equations of motion in the software environment MathCad 15.0, graphical dependencies for the main kinematic characteristics of the studied equipment are obtained (Fig. 4).

Theoretical analysis of the presented differential equations of motion of the executive equipment's of the developed vibrating disk crusher and graphical dependences (Fig. 4) are showed that during its operation without material supply, the resonant mode on the axis OZ to Az = 3.9 mm, on the axis OX-Ax = 2.25 mm at an angular velocity of the drive shaft 70 rad s⁻¹. As a result, the peak values of the total amplitude of oscillations are observed at 70 rad s⁻¹ and are A = 4.5 mm (Fig. 4a).

Analyzing the graphical dependence of the vibration speed on the angular velocity of the drive shaft operating modes are observed at values up to $0.3~{\rm m~s^{-1}}$ at 71 rad s⁻¹, peak values are observed in resonant mode at $150~{\rm s^{-1}}$ (Fig. 4b), and are $0.5~{\rm m~s^{-1}}$. Considering the graphical dependence of vibration acceleration (Fig. 4c) on angular velocity, the maximum value of $82~{\rm m~s^{2-1}}$ is observed at $150~{\rm s^{-1}}$.

Further analysis of the mathematical model in the use a wide range of analytics tools in the MathCad 15.0 software allowed to theoretically determine the operating frequency range of the vibrating disc crusher, in which the value of power consumption is close to the most rational values [6] and is N = 650... 750 W: $\omega = 100... 115$ rad / s, A = 3... 3.1 mm, v = 0.28 ... 0.31 m/s, a = 40... 43 m/s².

To verify the mathematical model, a series of experimental studies were performed and the real values of the amplitude-frequency characteristics of the developed equipment were established [20]. Experimental and theoretical graphs of the distribution of the main parameters of the studied system are presented (Fig. 5).

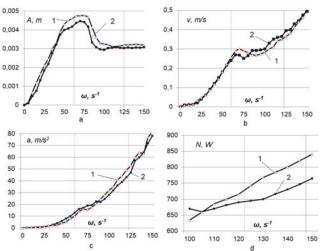


Fig. 5. Comparison of theoretical and experimental research results: a) amplitude of oscillations; b) vibration speed; c) vibration acceleration; d) power consumption; 1 – theoretical research; 2 – experimental research

Thus, it was found that the discrepancy between theoretical and experimental results is 7.2–12.1% and does not exceed the recommended value for vibratory crushers (up to 15%).

Conclusions

The result of the analysis of the proposed mathematical model of the vibrating crusher is the analytical and graphical dependences of the main kinematic parameters of oscillations. Recommended crusher operating parameters: ω =100...115 s⁻¹, A=3.1... 3.2 mm, v = 0.28 ... 0.31 m / s,

a=40...43 m/s². The energy consumption for the crusher drive is N=650...750 W. The peak values of the total amplitude in the resonant mode are A_{max} =4.5 mm at ω =70 s¹.

The comparative analysis of deviations of theoretical and experimental research on power and amplitude-frequency parameters of the developed equipment revealed a discrepancy of the received values within 7.2–12.1% that confirms the adequacy of the developed mathematical models. So, the proposed mathematical model with sufficient accuracy and reliability reflects the modes of oscillation of the machine and can be used when substantiating the parameters of the vibrating disc crusher.

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